A nonlocal convection-diffusion equation

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Few words about local diffusion problems

$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbb{R}^d \times (0, \infty), \\ u(0) = u_0. \end{cases}$$

For any $u_0 \in L^1(\mathbb{R})$ the solution $u \in C([0,\infty), L^1(\mathbb{R}^d))$ is given by:

$$u(t,x) = (H(t,\cdot) * u_0)(x)$$

where

$$H(t,x) = (4\pi t)^{d/2} \exp(-\frac{|x|^2}{4t})$$

Smoothing effect

 $u \in C^{\infty}((0,\infty), \mathbb{R}^d)$

Decay of solutions, $1 \le p \le q \le \infty$:

$$\|u(t)\|_{L^{q}(\mathbb{R}^{d})} \lesssim t^{-\frac{d}{2}(\frac{1}{p}-\frac{1}{q})} \|u_{0}\|_{L^{p}(\mathbb{R}^{d})}$$



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Zuazua& Duoandikoetxea, CRAS '92 For all $\varphi \in L^p(\mathbb{R}^d, 1+|x|^k)$

$$u(t,\cdot) \sim \sum_{|\alpha| \le k} \frac{(-1)^{|\alpha|}}{\alpha!} \left(\int f(x) x^{\alpha} dx \right) D^{\alpha} H(t,\cdot) \quad \text{in } L^q(\mathbb{R}^d)$$

for some $\boldsymbol{p},\boldsymbol{q},\boldsymbol{k}$



A linear nonlocal problem

E. Chasseigne, M. Chaves and J. D. Rossi, Asymptotic behavior for nonlocal diffusion equations, J. Math. Pures Appl., 86, 271–291, (2006).

$$\begin{cases} u_t(x,t) = J * u - u(x,t) = \int_{\mathbb{R}^d} J(x-y)u(y,t) \, dy - u(x,t), \\ = \int_{\mathbb{R}^d} J(x-y)(u(y,t) - u(x,t)) \, dy \\ u(x,0) = u_0(x), \end{cases}$$

where $J:\mathbb{R}^N\to\mathbb{R}$ be a nonnegative, radial function with $\int_{\mathbb{R}^N}J(r)dr=1$



P. Fife. Some nonclassical trends in parabolic and parabolic-like evolutions. Trends in nonlinear analysis, 153–191, Springer, Berlin, 2003.

- $u(\boldsymbol{x},t)$ the density of a single population at the point \boldsymbol{x} at time t
- $\bullet \ J(x-y)$ the probability distribution of jumping from y to x . Then

• $(J * u)(x, t) = \int_{\mathbb{R}^d} J(y - x)u(y, t) dy$ is the rate at which individuals are arriving to x from all other places and $-u(x, t) = -\int_{\mathbb{R}^d} J(y - x)u(x, t) dy$ is the rate at which they are leaving x to travel to all other sites.

Thus in the absence of external or internal sources, the density u satisfies the nonlocal equation (4).

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Heat equation and nonlocal diffusion

- Similarities
 - bounded stationary solutions are constant
 - a maximum principle holds for both of them

Difference

• there is no regularizing effect in general The fundamental solution can be decomposed as

$$w(x,t) = e^{-t}\delta_0(x) + v(x,t),$$
 (1)

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with v(x,t) smooth

 $S(t)\varphi = e^{-t}\varphi + v * \varphi =$ smooth as initial data + smooth part = no smoothing effect



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 $S(t)\varphi = e^{-t}\varphi + v * \varphi =$ smooth as initial data + smooth part = no smoothing effect • If $\hat{J}(\xi) = 1 - A|\xi|^{\alpha} + o(|\xi|^{\alpha}), \xi \sim 0$, the asymptotic behavior is the same as the one for solutions of the evolution given by the $\alpha/2$ fractional power of the laplacian:

$$\lim_{t \to +\infty} t^{d/\alpha} \max_{x} |u(x,t) - v(x,t)| = 0,$$

where v is the solution of $v_t(x,t) = -A(-\Delta)^{\alpha/2}v(x,t)$ with initial condition $v(x,0) = u_0(x)$.

• The asymptotic profile is given by

$$\lim_{t \to +\infty} \max_{y} \left| t^{d/\alpha} u(yt^{1/\alpha},t) - (\int_{\mathbb{R}^d} u_0) G_A(y) \right| = 0,$$

where $G_A(y)$ satisfies $\hat{G}_A(\xi) = e^{-A|\xi|^{\alpha}}$.



- I.L. Ignat and J.D. Rossi, *Refined asymptotic expansions for nonlocal diffusion equations*, submitted to Journal of Evolution Equations.
- I.L. Ignat and J.D. Rossi, Asymptotic behaviour for a nonlocal diffusion equation on a lattice, Accepted in Zeitschrift für Angewandte Mathematik und Physik (ZAMP).



The classical convection-diffusion equation

For $a \in \mathbb{R}^d$ and $q \geq 1$

$$\left\{\begin{array}{l} u_t - \Delta u = a \cdot \nabla(|u|^{q-1}u) \text{ in } (0,\infty) \times \mathbb{R}^d \\ u(0) = u_0 \end{array}\right.$$

• Asymptotic Behaviour by using

$$\frac{d}{dt} \int_{\mathbb{R}^d} |u|^p dx = -\frac{4(p-1)}{p} \int_{\mathbb{R}^d} |\nabla(|u|^{p/2})|^2 dx.$$

- M. Schonbek, *Uniform decay rates for parabolic conservation laws*, Nonlinear Anal., 10(9), 943–956, (1986).
- M. Escobedo and E. Zuazua, Large time behavior for convection-diffusion equations in R^N, J. Funct. Anal., 100(1), 119–161, (1991).

Methods to obtain asymptotics

• M. Schonbek \rightarrow Fourier Splitting Method Using that $\frac{d}{dt} \int_{\mathbb{R}^d} |u|^2 dx \leq -c \int_{\mathbb{R}^d} |\nabla u|^2 dx$, define $S(t) = \{\xi : |\xi| \leq (\frac{d}{c(t+1)})^{1/2}\}$. Then

$$\begin{aligned} \frac{d}{dt} \int u^2 &\leq -c \int_S |\xi|^2 \hat{u}^2 - c \int_{S^c} |\xi|^2 \hat{u}^2 \\ &\leq -\frac{d}{t+1} \int_{S^c} \hat{u}^2 = -\frac{d}{t+1} \int u^2 + \frac{d}{t+1} \int_S \hat{u}^2 \end{aligned}$$

$$\begin{split} \frac{d}{dt}(t+1)^d \int u^2 &\leq d(\frac{d}{t+1})^{d-1} \int_S \hat{u}^2 \leq c(d)(t+1)^{-d/2+1} \|\hat{u}\|_{\infty}^2 \\ &\leq c(d)(t+1)^{-d/2+1} \|u(t)\|_{L^1(\mathbb{R}^d)} \\ &\leq c(d)(t+1)^{-d/2+1} \|u_0\|_{L^1(\mathbb{R}^d)} \end{split}$$



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• M. Escobedo and E. Zuazua \rightarrow Energy method For $d \ge 3$ using the Sobolev inequality $||v||_{2d/(d-2)} \le C(d) ||\nabla v||_2$ with $v = |u|^{p/2}$ we get

$$||u||_{pd/(d-2)}^p \le C(d) ||\nabla(|u|^{p/2})||_2^2.$$

Interpolation inequality

$$||u||_p \le ||u||_{pd/(d-2)}^{d(p-1)/(2+d(p-1))} ||u||_1^{2/(2+d(p-1))}$$

gives us

$$\|u\|_p^{(d(p-1)+2)/d(p-1)} \le c(p,d)\|u\|_1^{2p/d(p-1)}\|\nabla(|u|^{p/2})\|_2^2.$$

Then

$$\frac{d}{dt}\|u(t)\|_{p}^{p} + \frac{C}{\|u_{0}\|_{1}^{2p/d(p-1)}}\|u(t)\|_{p}^{p[d(p-1)+2]/d(p-1)} \leq 0$$

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Nonlocal Convection-Diffusion

L.I. Ignat and J.D. Rossi, *A nonlocal convection-diffusion equation*, J. Funct. Anal., 251, 399–437, (2007).

$$egin{aligned} & (u_t(t,x) = (J * u - u) \, (t,x) + (G * (f(u)) - f(u)) \, (t,x), & t > 0, \, x \in U, \ & (u(0,x) = u_0(x), & x \in \mathbb{R}^d. \end{aligned}$$

- J and G are nonnegatives and verify $\int_{\mathbb{R}^d} J(x) dx = \int_{\mathbb{R}^d} G(x) dx = 1.$
- J smooth, $J \in S(\mathbb{R}^d)$, radially symmetric
- G smooth, G ∈ S(ℝ^d), but not necessarily symmetric, then the individuals have greater probability of jumping in one direction than in others, provoking a convective effect

•
$$\widehat{J}(\xi) - 1 + \xi^2 \sim |\xi|^3$$
, for ξ close to 0.

•
$$f(u) = |u|^{q-1}u$$
 with $q > 1$

Theorem

For any $u_0 \in L^1(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d)$ there exists a unique global solution

$$u \in C([0,\infty); L^1(\mathbb{R}^d)) \cap L^\infty([0,\infty); \mathbb{R}^d).$$

If u and v are solutions of (12) corresponding to initial data $u_0, v_0 \in L^1(\mathbb{R}^d) \cap L^{\infty}(\mathbb{R}^d)$ respectively, then the following contraction property

$$||u(t) - v(t)||_{L^1(\mathbb{R}^d)} \le ||u_0 - v_0||_{L^1(\mathbb{R}^d)}$$

holds for any $t \ge 0$. In addition,

$$\|u(t)\|_{L^{\infty}(\mathbb{R}^d)} \leq \|u_0\|_{L^{\infty}(\mathbb{R}^d)}.$$

Why not $u_0 \in L^1(\mathbb{R}^d)$?

- For the local problem $||v(t)||_{\infty} \leq C||v_0||_1 t^{-d/2}$ for any $v_0 \in L^1 \cap L^{\infty}$, so for any $u_0 \in L^1$ we can choose $u_{0,\epsilon} \in L^1 \cap L^{\infty}$ with $u_{0,\epsilon} \to u_0$ in L^1 and we can pass to the limit.
- In the nonlocal model, we cannot prove such type of inequality independently of the $L^\infty\text{-norm}$ of the initial data.
- In the one-dimensional case with $f(u) = |u|^{q-1}u$, $1 \le q < 2$ we have

$$\sup_{u_0 \in L^1(\mathbb{R})} \sup_{t \in [0,1]} \frac{t^{\frac{1}{2}} \|u(t)\|_{L^{\infty}(\mathbb{R})}}{\|u_0\|_{L^1(\mathbb{R})}} = \infty.$$

• The L^1-L^∞ regularizing effect is not available for the linear equation $w_t=J\ast w-w$

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Back to the linear semigroup

$$\begin{cases} w_t(t,x) = (J * w - w)(t,x), & t > 0, x \in \mathbb{R}^d, \\ w(0,x) = u_0(x), & x \in \mathbb{R}^d. \end{cases}$$
(2)

Lemma

The fundamental solution of (2) can be decomposed as

$$S(t,x) = e^{-t}\delta_0(x) + K_t(x),$$
 (3)

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with

$$K_t(x) = e^{-t} (\exp(t\widehat{J}(\xi)) - 1)^{\vee}$$

smooth.

Lemma

Let J be such that $\widehat{J}(\xi) - 1 + \xi^2 \sim |\xi|^3$, as $\xi \sim 0$. For any $p \ge 1$, K_t satisfies:

$$\|K_t\|_{L^p(\mathbb{R}^d)} \le c(p,J) \langle t \rangle^{-\frac{d}{2}(1-\frac{1}{p})}$$
(4)

for any t > 0.

Remark

Assumning that $J \in S(\mathbb{R}^d)$ and replacing t with $\langle t \rangle$, K_t has the same behaviour as the heat kernel H_t



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Back to the nonlinear problem: Local existence

Banach's Fix Point Theorem for the functional

$$\Phi[u](t) = S(t) * u_0 + \int_0^t S(t-s) * (G * (f(u)) - f(u))(s) \, ds$$

and the space

$$X(T) = C([0,T]; L^1(\mathbb{R}^d)) \cap L^\infty([0,T]; \mathbb{R}^d)$$

endowed with the norm

$$\|u\|_{X(T)} = \sup_{t \in [0,T]} \left(\|u(t)\|_{L^1(\mathbb{R}^d)} + \|u(t)\|_{L^\infty(\mathbb{R}^d)} \right).$$



For any
$$1 \le p \le \infty$$

$$\|S(t) * \varphi\|_p \le (e^{-t} + \|K_t\|_1) \|\varphi\|_p \le 3 \|\varphi\|_p$$

$$\|S(t) * G * \varphi - S(t) * \varphi\|_{L^p(\mathbb{R}^d)} \le C \langle t \rangle^{-\frac{d}{2}(\frac{1}{p} - \frac{1}{r}) - \frac{1}{2}} \|\varphi\|_{L^p(\mathbb{R}^d)}.$$
(5)



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Global existence: control of the L^1 and L^∞ norms

- L^1 -norm: easy, $rac{d}{dt}\int_{\mathbb{R}^d}|u(t,x)|\,dx\leq 0$
- L^{∞} -norm

$$\frac{d}{dt} \int_{\mathbb{R}^d} (u(t,x) - m)^+ dx = I_1(t) + I_2(t)$$

where

$$I_1(t) = \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} J(x-y)u(t,y)\operatorname{sgn}(u(t,x)-m)^+ dy dx$$
$$-\int_{\mathbb{R}^d} u(t,x)\operatorname{sgn}(u(t,x)-m)^+ dx$$

and

$$I_{2}(t) = \int_{\mathbb{R}^{d}} \int_{\mathbb{R}^{d}} G(x-y) f(u)(t,y) \operatorname{sgn}(u(t,x)-m)^{+} dy dx - \int_{\mathbb{R}^{d}} f(u)(t,x) \operatorname{sgn}(u(t,x)-m)^{+} dx.$$

Lemma

Let $\theta \in L^1(\mathbb{R}^d)$ and K be a nonnegative function with mass one. Then for any $\mu \geq 0$ the following hold:

$$\int_{\theta(x)>\mu} \int_{\mathbb{R}^d} K(x-y)\theta(y) \, dy \, dx \le \int_{\theta(x)>\mu} \theta(x) \, dx.$$
 (6)

Proof. Step 1. $\mu = 0 \rightarrow \mu > 0$. For $\mu = 0$

$$\begin{split} &\int_{\theta(x)>0} \int_{\mathbb{R}^d} K(x-y)\theta(y) \, dy \, dx \le \int_{\theta(x)>0} \int_{\theta(y)>0} K(x-y)\theta(y) \, dy \, dx \\ &= \int_{\theta(y)>0} \theta(y) \int_{\theta(x)>0} K(x-y) \, dx \, dy \le \int_{\theta(y)>0} \theta(y) \int_{\mathbb{R}^d} K(x-y) \, dx \, dy \\ &= \int_{\theta(y)>0} \theta(y) \, dy. \end{split}$$

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Due to the lack of regularizing effect, the $L^{\infty}(\mathbb{R})$ -norm does not get bounded for positive times when we consider initial conditions in $L^1(\mathbb{R})$. This is in contrast to what happens for the local convection-diffusion problem.

Proposition

Let
$$d = 1$$
 and $|f(u)| \le C|u|^q$ with $1 \le q < 2$. Then

$$\sup_{u_0 \in L^1(\mathbb{R})} \sup_{t \in [0,1]} \frac{t^{\frac{1}{2}} \|u(t)\|_{L^{\infty}(\mathbb{R})}}{\|u_0\|_{L^1(\mathbb{R})}} = \infty.$$



Assume the above limit to be finite. Then

$$\|u(1)\|_{L^{\infty}(\mathbb{R})} \ge \|S(1) * u_{0}\|_{L^{\infty}(\mathbb{R})} - \left\| \int_{0}^{1} S(1-s) * (G * (f(u)) - f(u))(s) \, ds \, \right\|_{L^{\infty}(\mathbb{R})}$$

and

$$\begin{split} \left\| \int_{0}^{1} S(1-s) * (G * (f(u)) - f(u))(s) \, ds \, \right\|_{L^{\infty}(\mathbb{R})} \\ & \leq \int_{0}^{1} \langle 1-s \rangle^{-\frac{1}{2}} \| f(u(s)) \|_{L^{\infty}(\mathbb{R})} \, ds \\ & \leq C \int_{0}^{1} \| u(s) \|_{L^{\infty}(\mathbb{R})}^{q} \, ds \leq C M^{q} \| u_{0} \|_{L^{1}(\mathbb{R})}^{q} \int_{0}^{1} s^{-\frac{q}{2}} \, ds \\ & \leq C M^{q} \| u_{0} \|_{L^{1}(\mathbb{R})}^{q}. \end{split}$$



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The $L^{\infty}(\mathbb{R})$ -norm of the solution at time t = 1 satisfies

$$\begin{aligned} \|u(1)\|_{L^{\infty}(\mathbb{R})} &\geq \|S(1) * u_{0}\|_{L^{\infty}(\mathbb{R})} - CM^{q} \|u_{0}\|_{L^{1}(\mathbb{R})}^{q} \\ &\geq e^{-1} \|u_{0}\|_{L^{\infty}(\mathbb{R})} - \|K_{1}\|_{L^{\infty}(\mathbb{R})} \|u_{0}\|_{L^{1}(\mathbb{R})} - CM^{q} \|u_{0}\|_{L^{1}(\mathbb{R})}^{q} \\ &\geq e^{-1} \|u_{0}\|_{L^{\infty}(\mathbb{R})} - C \|u_{0}\|_{L^{1}(\mathbb{R})} - CM^{q} \|u_{0}\|_{L^{1}(\mathbb{R})}^{q}. \end{aligned}$$

Choosing now a sequence $u_{0,\varepsilon}$ with $||u_{0,\varepsilon}||_{L^1(\mathbb{R})} = 1$ and $||u_{0,\varepsilon}||_{L^\infty(\mathbb{R})} \to \infty$ we obtain that

 $\|u_{0,\varepsilon}(1)\|_{L^{\infty}(\mathbb{R})} \to \infty.$



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Let us consider

$$J_{\varepsilon}(s) = \frac{1}{\varepsilon^d} J\left(\frac{s}{\varepsilon}\right), \qquad G_{\varepsilon}(s) = \frac{1}{\varepsilon^d} G\left(\frac{s}{\varepsilon}\right),$$

and the solution $u_{\varepsilon}(t,x)$ to our convection-diffusion problem rescaled adequately,

$$\begin{cases} (u_{\varepsilon})_{t}(t,x) = \frac{1}{\varepsilon^{2}} \int_{\mathbb{R}^{d}} J_{\varepsilon}(x-y)(u_{\varepsilon}(t,y) - u_{\varepsilon}(t,x)) \, dy \\ + \frac{1}{\varepsilon} \int_{\mathbb{R}^{d}} G_{\varepsilon}(x-y)(f(u_{\varepsilon}(t,y)) - f(u_{\varepsilon}(t,x))) \\ u_{\varepsilon}(x,0) = u_{0}(x). \end{cases}$$

$$(7)$$



Theorem

For any T > 0, we have

$$\lim_{\varepsilon \to 0} \sup_{t \in [0,T]} \|u_{\varepsilon} - v\|_{L^{2}(\mathbb{R}^{d})} = 0,$$

where v(t, x) is the unique solution to the local convection-diffusion problem $v_t(t, x) = \Delta v(t, x) + b \cdot \nabla f(v)(t, x)$ with initial condition $v(x, 0) = u_0(x) \in L^1(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$ and $b = (b_1, ..., b_d)$ given by

$$b_j = \int_{\mathbb{R}^d} x_j G(x) \, dx, \qquad j = 1, \dots, d.$$

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Proof of convergence

We write the two problems in the semigroup formulation,

$$u_{\varepsilon}(t) = S_{\varepsilon}(t) * u_{0} + \int_{0}^{t} \frac{S_{\varepsilon}(t-s) * G_{\varepsilon} - S_{\varepsilon}(t-s)}{\varepsilon} * f(u_{\varepsilon}(s)) ds$$

and

$$v(t) = H(t) * u_0 + \int_0^t b \cdot \nabla H(t-s) * f(v(s)) \, ds.$$

Then

$$\sup_{t\in[0,T]}\|u_{\varepsilon}(t)-v(t)\|_{L^2(\mathbb{R}^d)}\leq \sup_{t\in[0,T]}I_{1,\varepsilon}(t)+\sup_{t\in[0,T]}I_{2,\varepsilon}(t)$$

where $I_{1,\varepsilon}(t)$ and $I_{2,\varepsilon}(t)$ are given by the difference of linear, respectively nonlinear part.



Using the Fourier representation of the two semigroups and applying Lebesgue's Convergence Theorem we get

$$\|S_{\varepsilon}(t)*u_{0}-H(t)*u_{0}\|_{L^{2}(\mathbb{R}^{d})}^{2} = \int_{\mathbb{R}^{d}} \left|e^{t\frac{\widehat{i}(\varepsilon\xi)-1}{\varepsilon^{2}}} - e^{-t\xi^{2}}\right|^{2} |\widehat{u_{0}}(\xi)|^{2} d\xi \to 0$$



Lemma (Uniform decay with respect to ε)

There exists a positive constant C = C(J,G) such that

$$\left\| \left(\frac{S_{\varepsilon}(t) * G_{\varepsilon} - S_{\varepsilon}(t)}{\varepsilon} \right) * \varphi \right\|_{L^{2}(\mathbb{R}^{d})} \leq C t^{-\frac{1}{2}} \|\varphi\|_{L^{2}(\mathbb{R}^{d})}$$

holds for all t > 0 and $\varphi \in L^2(\mathbb{R}^d)$, uniformly on $\varepsilon > 0$.

Remark

The decay is the same as for ∇H :

$$\|\nabla H * \varphi\|_{L^2(\mathbb{R})} \le t^{-1/2} \|\varphi\|_{L^2(\mathbb{R})}.$$

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Lemma

Let be T > 0 and M > 0. Then the following

$$\lim_{\varepsilon \to 0} \sup_{t \in [0,T]} \int_0^t \left\| \left(\frac{S_{\varepsilon}(s) * G_{\varepsilon} - S_{\varepsilon}(s)}{\varepsilon} - b \cdot \nabla H(s) \right) * \varphi(s) \right\|_{L^2(\mathbb{R}^d)} ds = 0,$$

holds uniformly for all $\|\varphi\|_{L^{\infty}([0,T];L^{2}(\mathbb{R}^{d}))} \leq M$. Here $b = (b_{1},...,b_{d})$ is given by

$$b_j = \int_{\mathbb{R}^d} x_j G(x) \, dx, \qquad j = 1, ..., d.$$



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Theorem

Let $f = |u|^{q-1}u$ with q > 1 and $u_0 \in L^1(\mathbb{R}^d) \cap L^\infty(\mathbb{R}^d)$. Then, for every $p \in [1, \infty)$ the solution u of equation (12) verifies $\|u(t)\|_{L^p(\mathbb{R}^d)} \leq C(\|u_0\|_{L^1(\mathbb{R}^d)}, \|u_0\|_{L^\infty(\mathbb{R}^d)}) \langle t \rangle^{-\frac{d}{2}(1-\frac{1}{p})}.$ (8)



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Lemma

Let R and δ be such that the function \widehat{J} satisfies:

$$\widehat{J}(\xi) \le 1 - \frac{|\xi|^2}{2}, \ |\xi| \le R \text{ and } \widehat{J}(\xi) \le 1 - \delta, \ |\xi| \ge R.$$
 (9)

Let us assume that for any t > 0 the function $u : [0, \infty) \times \mathbb{R}^d \to \mathbb{R}$ satisfies

$$\frac{d}{dt}\int_{\mathbb{R}^d}|u(t,x)|^2\,dx\leq -\frac{c}{2}\int_{\mathbb{R}^d}\int_{\mathbb{R}^d}J(x-y)(u(t,x)-u(t,y))^2\,dx\,dy$$

Then there exists a constant $a = d/c\delta$ such that

$$\|u(at)\|_{L^{2}(\mathbb{R}^{d})} \leq \frac{\|u(0)\|_{L^{2}(\mathbb{R}^{d})}}{(t+1)^{\frac{d}{2}}} + \frac{(2\omega_{0})^{\frac{1}{2}}(2\delta)^{\frac{d}{4}}}{(t+1)^{\frac{d}{4}}}\|u\|_{L^{\infty}([0,\infty); L^{1}(\mathbb{R}^{d}))}.$$

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Why not an energy method

If we want to use energy estimates to get decay rates (for example in $L^2(\mathbb{R}^d)$), we arrive easily to

$$\frac{d}{dt} \int_{\mathbb{R}^d} |w(t,x)|^2 \, dx = -\frac{1}{2} \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} J(x-y) (w(t,x) - w(t,y))^2 \, dx \, dy$$

when we deal with a solution of the linear equation $w_t = J \ast w - w$ and to

$$\frac{d}{dt}\int_{\mathbb{R}^d}|u(t,x)|^2\,dx\leq -\frac{1}{2}\int_{\mathbb{R}^d}\int_{\mathbb{R}^d}J(x-y)(u(t,x)-u(t,y))^2\,dx\,dy$$

when we consider the convection-diffusion problem. However, we can not go further since an inequality of the form

$$\left(\int_{\mathbb{R}^d} |u(x)|^p \, dx\right)^{\frac{2}{p}} \leq C \int_{\mathbb{R}^d} \int_{\mathbb{R}^d} J(x-y)(u(x)-u(y))^2 \, dx \, dy$$

is not available for p > 2 (or we have no idea how to get something useful).

Asymptotic profile-Weakly Nonlinear Behaviour

Theorem

Let
$$f = |u|^{q-1}u$$
 with $q > (d+1)/d$ and $u_0 \in L^1(\mathbb{R}^d, 1+|x|) \cap L^{\infty}(\mathbb{R}^d)$. For any $p \in [2, \infty)$ the following holds

$$t^{-\frac{d}{2}(1-\frac{1}{p})} \| u(t) - H(t) \int_{\mathbb{R}^d} u_0(x) \, dx \|_{L^p(\mathbb{R}^d)} \le C(J, G, p, d) \, \alpha_q(t),$$

where

$$\alpha_q(t) = \begin{cases} \langle t \rangle^{-\frac{1}{2}} & \text{if } q \ge (d+2)/d, \\ \langle t \rangle^{\frac{1-d(q-1)}{2}} & \text{if } (d+1)/d < q < (d+2)/d. \end{cases}$$

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The proof consists in two ingredients

• Estimate the difference $S(t) * \varphi - MH(t)$, known from our previous paper

$$\begin{split} \|S(t) * \varphi - MH(t)\|_{L^{p}(\mathbb{R}^{d})} \leq & Ce^{-t} \|\varphi\|_{L^{p}(\mathbb{R}^{d})} \\ &+ C \|\varphi\|_{L^{1}(\mathbb{R}^{d}, |x|)} \langle t \rangle^{-\frac{d}{2}(1 - \frac{1}{p}) - \frac{1}{2}}, \, t > 0, \end{split}$$

• Show that the contribution from the nonlinear part goes to zero faster than the above difference.



- The second term in the asymptotic expansion of the nonlocal convection-diffusion equation
- For the linear part, to try to remove the convolution property of the function ${\cal J}$

$$u_t(t,x) = \int_{\mathbb{R}} J(x,y)(u(t,y)-u(t,x))dy$$

under suitable properties of J: symmetric, $\int_{\mathbb{R}}Jdx=\int_{\mathbb{R}}Jdy=$ 1, etc..._



For a proof we invite to a



- An energy method which will allow us to recover the decay properties
- Something similar to the viscous Bourger's equation $u_t u_{xx} = (u^2)_x$ where the asymptotic profile is given by the self-similar solutions
- Until now we treat the case q > 1 + 1/d, case when the diffusion play the important role. What it happens when q < 1 + 1/d? Is there something similar to the local case, where the convection part gives the asymptotic behaviour (Escobedo, Vazquez, Zuazua, ARMA 1993)



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