

FIO with quadratic complex phase,
a mathematical justification of
the semiclassical Herman-Kluk propagator

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collaboration with Torben SWART

- Semiclassical time-dependent Schrödinger equation

$$i\varepsilon \frac{\partial}{\partial t} \psi^\varepsilon = H^\varepsilon \psi^\varepsilon, \quad \psi^\varepsilon(0) = \psi_0^\varepsilon \in L^2(\mathbb{R}^d; \mathbb{C})$$

$$H^\varepsilon = Op_{\text{Weyl}}^\varepsilon h \text{ with } h(x, \xi) \text{ symbol in } \mathcal{C}^\infty(\mathbb{R}^{2d}; \mathbb{C})$$

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- Gaussian coherent state centered at $(q, p) \in \mathbb{R}^{2d}$

$$g_{q,p}^\varepsilon(x) = \frac{1}{(\pi\varepsilon)^{d/4}} e^{\frac{i}{\varepsilon} p \cdot (x-q)} e^{-\frac{|x-q|^2}{2\varepsilon}}$$

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Overcomplete basis of $L^2(\mathbb{R}^d; \mathbb{C})$

$$\text{Id}_{L^2} = \frac{1}{(2\pi\varepsilon)^d} \iint \left| \begin{array}{c} g_{q,p}^\varepsilon \end{array} \right\rangle \left\langle g_{q,p}^\varepsilon \right| \mathrm{d}q \mathrm{d}p$$

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Overcomplete basis of $L^2(\mathbb{R}^d; \mathbb{C})$

$$e^{\frac{i}{\varepsilon} t H^\varepsilon} = \frac{1}{(2\pi\varepsilon)^d} \iint \left| e^{\frac{i}{\varepsilon} t H^\varepsilon} g_{q,p}^\varepsilon \right\rangle \langle g_{q,p}^\varepsilon | \, dq dp$$

Semiclassical propagation of Gaussian coherent states

(Hagedorn '80)

$$e^{\frac{i}{\varepsilon}tH^\varepsilon}g_{q,p}^\varepsilon = e^{\frac{i}{\varepsilon}S^{\kappa^t}(q,p)}\tilde{g}_{\kappa^t(q,p)}^\varepsilon + O(\varepsilon)$$

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- $\kappa^t = (X^{\kappa^t}, \Xi^{\kappa^t})$ given by Hamilton equation of motion in phase-space

$$\frac{d}{dt}\kappa^t = \begin{pmatrix} 0 & \text{I} \\ -\text{I} & 0 \end{pmatrix} \nabla_{(x,\xi)} h \circ \kappa^t, \quad \kappa^0 = \text{Id}$$

- S^{κ^t} classical action of motion

$$S^{\kappa^t}(q,p) = \int_0^t \left[\frac{d}{d\tau} X^{\kappa^\tau}(q,p) \cdot \Xi^{\kappa^\tau}(q,p) - h \circ \kappa^\tau(q,p) \right] d\tau$$

Semiclassical approximation of the unitary propagator

- Thawed Gaussian Approximation (*Heller '75*)

$$U_{\text{TGA}}^{\varepsilon}(t) = \frac{1}{(2\pi\varepsilon)^d} \iint \left| e^{\frac{i}{\varepsilon} S^{\kappa^t}(q,p)} \tilde{g}_{\kappa^t}^{\varepsilon}(q,p) \right\rangle \langle g_{q,p}^{\varepsilon} | \, dq dp$$

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- Herman-Kluk Propagator (*Herman and Kluk '84*)

$$U_{\text{HK}}^{\varepsilon}(t) = \frac{1}{(2\pi\varepsilon)^d} \iint u^t(q,p) \left| e^{\frac{i}{\varepsilon} S^{\kappa^t}(q,p)} g_{\kappa^t}^{\varepsilon}(q,p) \right\rangle \langle g_{q,p}^{\varepsilon} | \, dq dp$$

Main result (Swart R. '07)

Assume that $h(x, \xi)$ is subquadratic *i.e.*

$$\|\partial_{(x,\xi)}^\alpha h\|_\infty < \infty, \quad |\alpha| \geq 2$$

then, for $-T \leq t \leq T$,

$$\left\| e^{\frac{i}{\varepsilon} t H^\varepsilon} - U_{\text{HK}}^\varepsilon(t) \right\|_{L^2 \rightarrow L^2} = O(\varepsilon^{N+1} |t|)$$

$$u_{HK}^t(q, p) = u_0^t(q, p) + \varepsilon u_1^t(q, p) + \cdots + \varepsilon^N u_N^t(q, p)$$

where u_k^t are solutions of transport equations with initial conditions

$$u_0^0 = 1, \quad u_k^0 = 0 \quad (1 \leq k \leq N)$$

Semiclassical FIO with complex quadratic phase associated to a canonical transformation

$$[I^\varepsilon(\kappa; u)\varphi](x) = \frac{2^{d/2}}{(2\pi\varepsilon)^{3d/2}} \iiint e^{\frac{i}{\varepsilon}\Phi^\kappa(x,y,q,p)} u(q,p)\varphi(y)dydqdp$$

$$\text{with } \Phi^\kappa(x, y, q, p) = S^\kappa + \Xi^\kappa \cdot (x - X^\kappa) - p \cdot (y - q) \\ + \frac{i}{2}|x - X^\kappa|^2 + \frac{i}{2}|y - q|^2$$

$$\partial_q S^\kappa = -p + (\partial_q X^\kappa)\Xi^\kappa, \quad \partial_p S^\kappa = (\partial_p X^\kappa)\Xi^\kappa$$

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$$\text{for symbol } u, \quad I^\varepsilon(\kappa; u) : \mathcal{S}(\mathbb{R}^d; \mathbb{C}) \rightarrow \mathcal{S}(\mathbb{R}^d; \mathbb{C})$$

Sketch of the proof

$$\left[i\varepsilon \frac{\partial}{\partial t} - H^\varepsilon \right] I^\varepsilon \left(\kappa^t; \sum_{k=0}^N \varepsilon^k u_k^t \right) =$$
$$I^\varepsilon \left(\kappa^t; \sum_{k=0}^N \varepsilon^k v_k^t + \varepsilon^{N+1} v_{[N+1,}^{t,\varepsilon} \right)$$

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as consequence of composition theorems (*Swart R. '07*)

$$[Op_{\text{Weyl}}^\varepsilon h] \circ I^\varepsilon(\kappa; u) = I^\varepsilon \left(\kappa; \sum_{k=0}^N \varepsilon^k v_k + \varepsilon^{N+1} v_{[N+1]}^\varepsilon \right)$$

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where $v_0(q, p) = (h \circ \kappa)u$, $v_1 = \dots$

$\dots$ and $v_{[N+1]}^\varepsilon(x, y, q, p)$

- “eikonal” equation

$$v_0^t = \left[-\frac{d}{dt} S^{\kappa^t} + \frac{d}{dt} X^{\kappa^t} . \Xi^{\kappa^t} + h \circ \kappa^t \right] u_0^t$$

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- transport equations

$$v_1^t = i \frac{d}{dt} u_0^t - \frac{i}{2} \text{Tr} \left[\left(\mathcal{Z}^{\kappa^t} \right)^{-1} \frac{d}{dt} \mathcal{Z}^{\kappa^t} \right] u_0^t$$

$$\mathcal{Z}^{\kappa^t} := (\partial_p + i\partial_q)(\Xi^{\kappa^t} - iX^{\kappa^t})$$

Hierarchy of equations

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Herman-Kluk prefactor

$$u_0^t = \left(\det \mathcal{Z}^{\kappa^t} \right)^{1/2}$$

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$$v_{k+1}^t = i \frac{d}{dt} u_k^t - \frac{i}{2} \text{Tr} \left[\left(\mathcal{Z}^{\kappa^t} \right)^{-1} \frac{d}{dt} \mathcal{Z}^{\kappa^t} \right] u_k^t + L_k^t(u_0^t, \dots, u_{k-1}^t)$$

Final step of the proof

$$\left[i\varepsilon \frac{\partial}{\partial t} - H^\varepsilon \right] I^\varepsilon \left(\kappa^t; \sum_{k=0}^N \varepsilon^k u_k^t \right) = \varepsilon^{N+1} I^\varepsilon \left(\kappa^t; v_{[N+1]}^{t,\varepsilon} \right)$$

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We combine Gronwall's lemma with
 L^2 -boundedness theorems (*Swart R. '07*)

$$\|I^\varepsilon(\kappa; u)\|_{L^2 \rightarrow L^2} \leq C \sum_{|\alpha| \leq 2d+1} \|\partial_{(x,y)}^\alpha u(x, y, q, p)\|_\infty$$