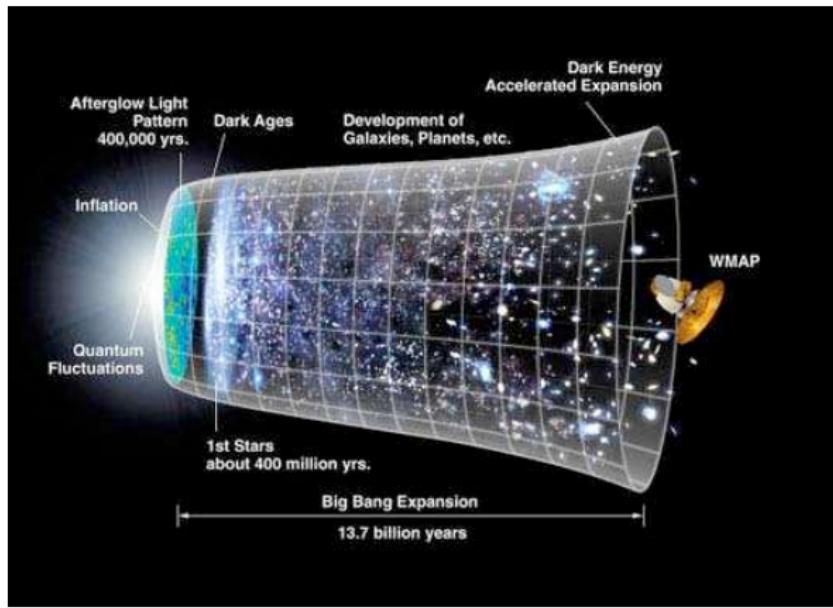


A GLIMPSE OF NONCOMMUTATIVE CURVATURE

Henri Moscovici

8 September 2010
IMAR, Bucharest



		I	II	III	
		Quarks			Force Carriers
Leptons	u	c	t	γ	
	d	s	b	g	
	ν_e	ν_μ	ν_τ	Z	
	e	μ	τ	W	

Three Generations of Matter

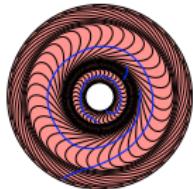
Model for noncommutative space: leaf space

Complete transversal

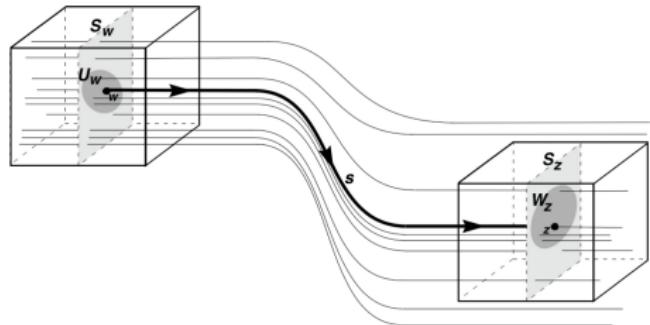


Foliation

Holonomy



Transversals



10

Spectral model: Dirac operator

- **Spin^c manifold:** closed Riemannian manifold (M^{2n}, g) such that the Clifford algebra bundle has “square root”:

$$\wedge^\bullet T^* M \otimes \mathbb{C} \cong \text{Cliff}(M) \otimes \mathbb{C} \cong \text{End}(\mathcal{S}) \cong \mathcal{S}^* \otimes \mathcal{S}$$

- **Spin^c-Dirac operator:** “square root” of Laplace operator:

$$\not{D} = \not{D}^*, \quad \not{D}^2 = \Delta^{\mathcal{S}} + \frac{1}{4}\kappa_g \quad (\text{scalar curvature})$$

- **Dirac spectral triple:** $(C^\infty(M^{2n}), L^2(\mathcal{S}^\pm), \not{D}^\pm)$,

$$\not{D} = \not{D}^*, \quad \not{D} \circ \gamma = -\gamma \circ \not{D}, \quad \gamma = \text{chirality operator}$$

$$(\not{D}^2 + 1)^{-\frac{1}{2}} \in \mathcal{L}^{2n+\varepsilon},$$

$$[\not{D}, f] := c(df) \in \mathcal{L}(\mathfrak{H}), \quad \forall f \in \mathcal{A} = C^\infty(M).$$

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Spectral presentation of spaces

- **Spectral triple:** $(\mathcal{A}, \mathfrak{H}, D)$, $\mathcal{A} \subset \mathcal{L}(\mathfrak{H})$ involutive subalgebra

$$D = D^* \quad \text{unbounded} \quad (F = \text{Sign } D)$$

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- **Simplest example: the circle**

$$u^{-1}[D, u] = 1, \quad uu^* = u^*u = 1,$$

$$\leadsto S^1 = \text{Spec}(\mathcal{A}), \quad \mathcal{A} = C^\infty(u), \quad D = -\frac{1}{i} \frac{d}{dx}$$

$$\Rightarrow d(x_1, x_2) = \sup_{\|[D, f]\| \leq 1} |x_1(f) - x_2(f)|.$$

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Local-Global Principle in Global Analysis

Theorem (ATIYAH-SINGER)

M^{2n} = closed spin manifold, D^\pm = Dirac operator, E = vector bundle

$$\text{Index}(D_E^+) = \int_M \hat{A}(R_g) ch(E, \nabla)$$

- cohomological expression to the index pairing:

$$\begin{array}{ccccc} K_*(M) & \otimes & K^*(M) & \xrightarrow{\text{Index}} & \mathbb{Z} \\ \downarrow ch_* & & \downarrow ch^* & & \downarrow \\ H_*(M) & \otimes & H^*(M) & \longrightarrow & \mathbb{R} \end{array}$$

- $\lim_{t \nearrow \infty} \text{Tr}(\gamma e^{-t D_E^2}) = \lim_{t \searrow 0} \text{Tr}(\gamma e^{-t D_E^2})$

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Cyclic cohomology and index pairing [Connes (1981)]

Let A be a C^* -algebra represented in $\mathcal{L}(\mathfrak{H})$, and let

$$F \in \mathcal{L}(\mathfrak{H}), \quad F^2 = I, \quad F\gamma = -\gamma F, \quad [F, a] \in \mathcal{L}^p(\mathfrak{H}), \quad \forall a \in \mathcal{A} \subset A.$$

CONNES turned the 'parametrix' index formula

$$\text{Index}(e F^+ e) = (-1)^\ell \text{Tr}(\gamma e [F, e]^n), \quad n = 2\ell > p,$$

into a **cyclic cohomological formulation of the index pairing**:

$$K^*(A) \otimes K_*(A) \rightarrow HC^*(\mathcal{A}) \otimes HC_*(\mathcal{A}) \rightarrow \mathbb{Z}$$

$$\text{Index}(e F^+ e) =: \langle [F], [e] \rangle = \langle ch^*(F), ch_*(e) \rangle$$

where

- $ch^*(F) := [\tau_F^n] \in HC^0(\mathcal{A})$, with $\tau_F^n \in HC_\lambda^n(\mathcal{A})$ given by

$$\tau_F^n(a^0, a^1, \dots, a^n) := c_n \text{Tr}(\gamma F[F, a^0] \dots [F, a^n]), \quad a^i \in \mathcal{A};$$

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Locality and zeta-function regularization

- Spectral triple $(\mathcal{A}, \mathfrak{H}, D)$ with **discrete dimension spectrum**: \exists discrete subset $\Sigma \subset \mathbb{C}$, such that $\zeta_b(z) = \text{Tr}(b|D|^{-z})$, $\Re z > p$, $b \in \mathcal{B}$, **admit holomorphic extensions to $\mathbb{C} \setminus \Sigma$** , and $\Gamma(z)\zeta_b(z)$ decays rapidly on finite vertical strips.

Simple dimension spectrum: Σ consists of simple poles.

- Then the residue functional

$$\int P := \text{Res}_{z=0} \zeta_P(2z), \quad P \in \Psi^*(\mathcal{A}, \mathfrak{H}, D)$$

is an algebraic trace (assuming s. d. s.).

- $\int P$ vanishes on $\Psi^*(\mathcal{A}, \mathfrak{H}, D) \cap \mathcal{L}^1(\mathfrak{H})$, and thus descends to a trace on the algebra of symbols $\mathcal{CS}^*(\mathcal{A}, \mathfrak{H}, D)$.

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Local Index Formula in NcG [Connes - M, 1995]

For an (odd) spectral triple $(\mathcal{A}, \mathfrak{H}, D)$ such that \exists residues

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- ① $[(\varphi_n)_{n=1,3,\dots}]$ is a cocycle in the (b, B) -bicomplex of \mathcal{A} ,

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Example: the Dirac spectral triple

- ① The dimension spectrum of the Dirac spectral triple $(C^\infty(M^m), L^2(\mathcal{S}), \not{D})$ is simple, and $= \{1, \dots, m\}$
- ② $\int P = \text{Wodzicki-Guillemain residue } (P), \quad \forall P \in \Psi DO(M)$
- ③ $\int f^0 [\not{D}, f^1]^{(k_1)} \dots [\not{D}, f^n]^{(k_n)} |\not{D}|^{-(n+2|k|)} = 0, \quad \text{if } |k| > 0$
- ④ $\int f^0 [\not{D}, f^1] \dots [\not{D}, f^n] |\not{D}|^{-n} = c_n \int_M \det \left(\frac{\nabla^2 / 4\pi i}{\sinh \nabla^2 / 4\pi i} \right)^{\frac{1}{2}} \wedge f^0 df^1 \wedge \dots \wedge df^n$
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Twisted spectral triples

- **Twisted commutators** : $(\mathcal{A}, \mathfrak{H}, D, \sigma)$, where $\sigma \in \text{Aut}(\mathcal{A})$,

$$[D, a]_\sigma := D a - \sigma(a) D$$

- Example: codimension 1 foliation:

$\mathcal{A}_G = C^\infty(S^1) \rtimes G$, $f U_\phi^{-1} \cdot g U_\psi^{-1} := f \phi^*(g) U_{\phi\psi}^{-1}$,
acting on $\mathfrak{H} = L^2(S^1) := \{\xi = \xi(x) dx^{\frac{1}{2}} \mid \xi \in L^2\}$, via

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Conformal twisting

- **Dirac spectral triple** Let $g' = e^{-4h} g$. Natural identifications of spinorial frames [*Hitchin, Bourguignon-Goduchon*] give to the transported Dirac operator the expression

$${}^g \not{D}^{g'} = e^{(n+1)h} \circ \not{D}^g \circ e^{-(n-1)h};$$

when taking into account the change $\text{vol}_{g'} = e^{-2nh} \text{vol}_g$,

$$\implies {}^g \not{D}^{g'} = e^h \not{D}^g e^h.$$

- **General spectral triple.** Given $(\mathcal{A}, \mathfrak{H}, D)$ and $h = h^* \in \mathcal{A}$ define $D_h = e^h D e^h$. $[D_h, a]$ is **unbounded** if h is not central. However, with the twisting automorphism

$$\sigma(a) = e^{2h} a e^{-2h} \implies [D_h, a]_\sigma := D_h a - \sigma(a) D_h \in \mathcal{L}(\mathfrak{H}).$$

Conformal twisting

- **Dirac spectral triple** Let $g' = e^{-4h} g$. Natural identifications of spinorial frames [*Hitchin, Bourguignon-Goduchon*] give to the transported Dirac operator the expression

$${}^g \not{D}^{g'} = e^{(n+1)h} \circ \not{D}^g \circ e^{-(n-1)h};$$

when taking into account the change $\text{vol}_{g'} = e^{-2nh} \text{vol}_g$,

$$\implies {}^g \not{D}^{g'} = e^h \not{D}^g e^h.$$

- **General spectral triple.** Given $(\mathcal{A}, \mathfrak{H}, D)$ and $h = h^* \in \mathcal{A}$ define $D_h = e^h D e^h$. $[D_h, a]$ is **unbounded** if h is not central.
However, with the twisting automorphism

$$\sigma(a) = e^{2h} a e^{-2h} \implies [D_h, a]_\sigma := D_h a - \sigma(a) D_h \in \mathcal{L}(\mathfrak{H}).$$

Noncommutative 2-torus

- $A_\theta := C^*\{U, V\}$, U, V unitaries, $VU = e^{2\pi i \theta} UV$, $\theta \in \mathbb{R} \setminus \mathbb{Q}$.
- Action of $s \in \mathbb{R}^2$, $\alpha_s(U^m V^n) = e^{i \langle s, (m, n) \rangle} U^m V^n$.
- Smooth subalgebra:
$$A_\theta^\infty = \left\{ \sum_{m, n \in \mathbb{Z}} a_{m, n} U^m V^n : |m|^k |n|^q |a_{m, n}| \leq C_{k, q}, \quad \forall k, q > 0 \right\}.$$
- Infinitesimal generators : derivations $\delta_1, \delta_2 : A_\theta^\infty \rightarrow A_\theta^\infty$,
$$\delta_1(U) = U, \quad \delta_1(V) = 0, \quad \delta_2(U) = 0, \quad \delta_2(V) = V.$$
- Unique normalized trace: $\tau(\sum_{m, n \in \mathbb{Z}} a_{m, n} U^m V^n) = a_{0,0}$.
- Dirac operator on $\mathcal{H}_0 = \bar{A}_\theta$, w.r.t. $\langle a, b \rangle = \tau(b^* a)$,

$$D_z = \begin{pmatrix} 0 & \delta_1 + z\delta_2 \\ -\delta_1 - \bar{z}\delta_2 & 0 \end{pmatrix}, \quad \Im z > 0.$$

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Conformal change on A_θ

- **Conformal state:** $\varphi(a) = \tau(ae^{-h})$,

$$\varphi(ab) = \varphi(b\sigma_i(a)), \quad a, b \in A_\theta,$$

i.e. φ is KMS with 1-parameter group

$$\sigma_t(x) = e^{ith}xe^{-ith} = \Delta^{-it}, \quad t \in \mathbb{R}.$$

- **Modular operator:** $\Delta(x) = e^{-h}xe^h$; σ_t generated by the derivation $-\log \Delta(x) = [h, x]$, $x \in A_\theta^\infty$.
- $\mathcal{H}_\varphi^+ :=$ Hilbert space completion of A_θ w.r.t. $\langle a, b \rangle_\varphi = \varphi(b^*a)$.
- **Bi-module of $(1,0)$ -forms** $\mathcal{H}^{(1,0)} =$ Hilbert space completion of the space of $\sum a \partial b$, $a, b \in A_\theta^\infty$ w.r.t.
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Associated spectral triples over A_θ

- $\partial_z = \delta_1 + z\delta_2, \quad \partial_z^* = \delta_1 + \bar{z}\delta_2, \quad z = x + iy, \quad y > 0,$

$$\mathcal{H}_\varphi = \mathcal{H}_\varphi^+ \oplus \mathcal{H}^{(1,0)}, \quad D_\varphi = \begin{pmatrix} 0 & \partial_\varphi^* \\ \partial_\varphi & 0 \end{pmatrix}, \quad \partial_\varphi = \partial : \mathcal{H}_\varphi^+ \rightarrow \mathcal{H}^{(1,0)};$$

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- $\zeta_\varphi(s) := \text{Tr}(|D_\varphi|^{-s}) = 2\text{Tr}((e^{\frac{h}{2}} \Delta e^{\frac{h}{2}})^{-s/2}).$

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Gauss-Bonnet

For a Riemannian surface Σ , $\zeta(s) = \text{Tr}(\Delta^{-s})$ satisfies

$$\zeta(0) + \dim \text{Ker} \Delta = \frac{1}{12\pi} \int_{\Sigma} K = \frac{1}{6} \chi(\Sigma),$$

in particular, $\zeta(0)$ is a topological invariant.

Theorem (Connes-Tretkoff)

For $\theta \in \mathbb{R} \setminus \mathbb{Q}$ and $z = 0$, let $h \in A_{\theta}^{\infty}$, and $\Delta_{\varphi} = \partial_{\varphi}^* \partial_{\varphi} \sim e^{\frac{h}{2}} \Delta e^{\frac{h}{2}}$. Then $\zeta_{\Delta_{\varphi}}(0)$ of the operator is independent of h .

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Pseudodifferential calculus [Connes,1980]

Symbols of order n : \mathcal{S}_n formed of smooth maps $\rho : \mathbb{R}^2 \rightarrow A_\theta^\infty$ such that
 $\forall i_1, i_2, j_1, j_2 \in \mathbb{Z}^+, \exists$ constant $C_{i,j}$,

$$\|\delta_1^{i_1} \delta_2^{i_2} \partial_1^{j_1} \partial_2^{j_2} \rho(\xi)\| \leq C_{i,j} (1 + |\xi|)^{n-j_1-j_2},$$

and $\exists \sigma \in C^\infty(\mathbb{R}^2 \setminus \{0\}, A_\theta^\infty)$ such that

$$\lim_{\lambda \rightarrow \infty} \lambda^{-n} \rho(\lambda \xi_1, \lambda \xi_2) = \sigma(\xi_1, \xi_2).$$

Associated pseudodifferential operator on A_θ^∞ :

$$P_\rho(a) = (2\pi)^{-2} \iint e^{-is \cdot \xi} \rho(\xi) \alpha_s(a) ds d\xi.$$

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Symbolic calculations

$$\sigma(k \Delta k + 1) = a_2(\xi) + a_1(\xi) + a_0(\xi), \quad k = e^{\frac{h}{2}}$$

where

$$a_2(\xi) = \xi_1^2 k^2 + |z|^2 \xi_2^2 k^2 + 2x \xi_1 \xi_2 k^2,$$

$$a_1(\xi) = 2\xi_1 k \delta_1(k) + 2|z|^2 \xi_2 k \delta_2(k) + 2x \xi_1 k \delta_2(k) + 2x \xi_2 k \delta_1(k),$$

$$a_0(\xi) = k \delta_1^2(k) + |z|^2 k \delta_2^2(k) + 2xk \delta_1 \delta_2(k).$$

Inverse symbol = $b_0 + b_1 + b_2 + \dots$, with

$$b_n = - \sum_{\substack{2+j+\ell_1+\ell_2-i=n, \\ 0 \leq j < n, 0 \leq i \leq 2}} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2}(b_j) \delta_1^{\ell_1} \delta_2^{\ell_2}(a_i) b_0,$$

of order $-2 - n$.

Symbolic calculations

$$\sigma(k \Delta k + 1) = a_2(\xi) + a_1(\xi) + a_0(\xi), \quad k = e^{\frac{h}{2}}$$

where

$$a_2(\xi) = \xi_1^2 k^2 + |z|^2 \xi_2^2 k^2 + 2x \xi_1 \xi_2 k^2,$$

$$a_1(\xi) = 2\xi_1 k \delta_1(k) + 2|z|^2 \xi_2 k \delta_2(k) + 2x \xi_1 k \delta_2(k) + 2x \xi_2 k \delta_1(k),$$

$$a_0(\xi) = k \delta_1^2(k) + |z|^2 k \delta_2^2(k) + 2xk \delta_1 \delta_2(k).$$

Inverse symbol = $b_0 + b_1 + b_2 + \dots$, with

$$b_n = - \sum_{\substack{2+j+\ell_1+\ell_2-i=n, \\ 0 \leq j < n, 0 \leq i \leq 2}} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2} (b_j) \delta_1^{\ell_1} \delta_2^{\ell_2} (a_i) b_0,$$

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$$\textcolor{red}{b_0} = (\xi_1^2 k^2 + |z|^2 \xi_2^2 k^2 + 2x \xi_1 \xi_2 k^2 + 1)^{-1},$$

$$\textcolor{red}{b_1} = -(b_0 a_1 b_0 + b_1 a_1 b_0 + \partial_1(b_0) \delta_1(a_2) b_0 + \partial_2(b_0) \delta_2(a_2) b_0),$$

$$\begin{aligned}\textcolor{blue}{b_2} = & -b_0 k \delta_1^2(k) b_0 - 2x b_0 k \delta_1 \delta_2(k) b_0 - |z|^2 b_0 k \delta_2^2(k) b_0 + \\& 6\xi_1^2 b_0^2 k^2 \delta_1(k)^2 b_0 + \xi_1^2 b_0^2 k^2 \delta_1^2(k) b_0 k + 5\xi_1^2 b_0^2 k^3 \delta_1^2(k) b_0 + \\& 2\xi_1^2 b_0 k \delta_1(k) b_0 \delta_1(k) b_0 k + 6\xi_1^2 b_0 k \delta_1(k) b_0 k \delta_1(k) b_0 + \\& 6x \xi_1^2 b_0^2 k^2 \delta_1(k) \delta_2(k) b_0 + 6x \xi_1^2 b_0^2 k^2 \delta_2(k) \delta_1(k) b_0 + \\& 2x \xi_1^2 b_0^2 k^2 \delta_1 \delta_2(k) b_0 k + 10x \xi_1^2 b_0^2 k^3 \delta_1 \delta_2(k) b_0 + \\& 2x \xi_1^2 b_0 k \delta_1(k) b_0 \delta_2(k) b_0 k + 6x \xi_1^2 b_0 k \delta_1(k) b_0 k \delta_2(k) b_0 + \\& 2x \xi_1^2 b_0 k \delta_2(k) b_0 \delta_1(k) b_0 k + 6x \xi_1^2 b_0 k \delta_2(k) b_0 k \delta_1(k) b_0 + \\& 12x \xi_1 \xi_2 b_0^2 k^2 \delta_1(k)^2 b_0 + 2x \xi_1 \xi_2 b_0^2 k^2 \delta_1^2(k) b_0 k + \\& 2x \xi_1 \xi_2 b_0^2 k^2 \delta_1^2(k) b_0 k 10x \xi_1 \xi_2 b_0^2 k^3 \delta_1^2(k) b_0 + \dots\end{aligned}$$

(... + 100s more such terms)

Scalar curvature

Theorem (Connes)

Let θ be an irrational number and k an invertible positive element of A_θ^∞ . Then the value at the origin of the zeta function of the twisted spectral triple $(A_\theta^\infty, \mathcal{H}, D)$ is given for any $a \in A_\theta^\infty$ by

$$\mathrm{Tr}(a|D|^{-s})|_{s=0} = \tau(a \left(K(\log \Delta)(\Delta(h)) + \sum_{j=1}^2 K(\nabla_{(1)}, \nabla_{(2)}) \delta_j(h) \delta_j(h) \right)).$$

Thus, the scalar curvature of the metric D_φ with $\Delta_\varphi \cong e^{\frac{h}{2}} \Delta e^{\frac{h}{2}}$ is

$$\kappa_h = K(\log \Delta)(\Delta(h)) + \sum_{j=1}^2 K(\nabla_{(1)}, \nabla_{(2)}) \delta_j(h) \delta_j(h),$$

where $\nabla_{(i)}$ is the action of $\log \Delta$ on the i -th factor of the product.

Scalar curvature functions

$$K(u) = \frac{e^u - e^{u/2}u - 1}{(-1 + e^{u/2})^3}, \quad K(s, t) = \left(1 + \operatorname{ch}\left(\frac{s+t}{2}\right)\right) K_0(s, t)$$

$$K_0(s, t) =$$

$$\frac{-t(s+t)\operatorname{ch}(s) + s(s+t)\operatorname{ch}(t) - (s-t)(s+t + \operatorname{sh}(s) + \operatorname{sh}(t) - \operatorname{sh}(s+t))}{st(s+t)\operatorname{sh}\left(\frac{s}{2}\right)\operatorname{sh}\left(\frac{t}{2}\right)\operatorname{sh}\left(\frac{s+t}{2}\right)^2}$$

$$K_0(t, s) = -K_0(s, t); \quad K_0(-s, -t) = -K_0(s, t).$$

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Conformal index of a spectral triple

$(\mathcal{A}, \mathfrak{H}, D)$ = *p*-summable spectral triple with good pseudodifferential calculus, i.e. \exists asymptotic expansions

$$\mathrm{Tr} \left(A e^{-tD_s^2} \right) =_{t \searrow 0} \sum_{j=0}^{\infty} a_j(A, s) t^{\frac{j-N-p}{2}} + 0(1),$$

for any $A \in \mathcal{D}^N(\mathcal{A}, \mathfrak{H}, D)$; $h = h^* \in \mathcal{A}$, $D_s = e^{sh} D e^{sh}$.

Moreover, assume good resolvent approximation, implying that these expansions can be differentiated in the term-by-term in $s \in [-1, 1]$.

Theorem (Conformal index à la Branson-Ørsted)

$\zeta_D(0)$ is a conformal invariant.

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Proof.

$$D_s = e^{sh} D e^{-sh} \implies \frac{de^{-D_s^2}}{ds} = - \int_0^1 e^{-uD_s^2} \frac{dD_s^2}{ds} e^{-(1-u)D_s^2} du$$

$$\frac{dD_s}{ds} \equiv \frac{d}{ds}(e^{sh} D e^{-sh}) = hD_s + D_sh,$$

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Equating the coefficients of the asymptotic expansions,

$$\frac{d}{ds} a_j(s) = 2(j-p) a_j(h, s), \quad \text{in particular} \quad \frac{d}{ds} a_p(s) = 0.$$

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Corollary (1) (Connes-Tretkoff, Khalkhali-Fathi Zadeh)

Gauss-Bonnet for the noncommutative 2-torus with any translation invariant metric.

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Same argument as above, but for the Laplacian $\Delta_s = e^{sh} \Delta e^{sh}$:

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□

Corollary (2) (Connes-Chamseddine)

The constant term in the SM spectral action $\text{Tr } F \left(\frac{D^2}{m^2} \right) + \langle \Psi | D | \Psi \rangle$ is independent of dilaton field rescaling $D^2 \mapsto e^h D^2 e^h$.

$$\text{Tr } F \left(\frac{D^2}{m^2} \right) \simeq \frac{m^4}{2} f_0 \int D^{-4} + \frac{m^2}{2} f_2 \int D^{-2} + f_4 \zeta_{D^2}(0) + \dots$$

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